

Slater-Pauling behavior and origin of the half-metallicity of the full-Heusler alloys

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Using the full-potential screened Korringa-Kohn-Rostoker method we study the full-Heusler alloys based on Co, Fe, Rh, and Ru. We show that many of these compounds show a half-metallic behavior; however, in contrast to the half-Heusler alloys the energy gap in the minority band is extremely small due to states localized only at the Co (Fe, Rh, or Ru) sites which are not present in the half-Heusler compounds. The full-Heusler alloys show a Slater-Pauling behavior and the total spin magnetic moment per unit cell (M_t) scales with the total number of valence electrons (Z_t) following the rule $M_t = Z_t - 24$. We explain why the spin-down band contains exactly 12 electrons using arguments based on group theory and show that this rule holds also for compounds with less than 24 valence electrons. Finally we discuss the deviations from this rule and the differences compared to the half-Heusler alloys.

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I. INTRODUCTION

Increased interest in the field of magnetoelectronics or spin electronics during the last decade¹ has intensified research on the so-called half-ferromagnetic materials. The latter ones present a gap in the minority band and thus can be used as perfect spin filters or to enhance the performance of spin-dependent devices as electrons at the Fermi level are 100% spin polarized. The first material which was predicted to be a half-ferromagnet was the half-Heusler alloy NiMnSb found by de Groot and collaborators² in 1983. This prediction has been verified also by other authors³⁻⁵ and the half-ferromagnetic character has been also well established experimentally both by using positron annihilation experiments⁶ or inverse photoemission.⁷ Recently there has been an increased interest in thin films of this material both experimentally⁸ and using first-principles calculations.^{9,10}

Although half-Heusler alloys like NiMnSb have attracted a lot of interest, the second family of Heusler compounds, the so-called full-Heusler alloys, have been studied much more extensively due to the existence of diverse magnetic phenomena,^{11,12} mainly the transition from a ferromagnetic phase to an antiferromagnetic one by changing the concentration of the carriers.¹³ The full-Heusler alloys have the type X_2YZ (see Fig. 1) and they crystallize in the $L2_1$ structure which consists of four fcc sublattices. Ziebeck and Webster¹⁴ were the first to synthesize full-Heusler alloys containing Co, and Ishida and collaborators^{15,16} have proposed that compounds of the type Co_2MnZ , where Z stands for Si and Ge, are also half-ferromagnets. Also Heusler alloys of the type Fe_2MnZ have been proposed to show half-ferromagnetism.¹⁷ But Brown *et al.*¹⁸ using polarized neutron diffraction measurements have shown that there is a finite very small spin-down density of states (DOS) at the Fermi level instead of an absolute gap in agreement with the *ab initio* calculations of Kübler *et al.* for Co_2MnAl and Co_2MnSn compounds.¹³ Recently, Ambrose *et al.*¹⁹ managed to grow a Co_2MnGe thin

film on a GaAs(001) substrate by molecular beam epitaxy and have proved the creation of domains during the growth.²⁰ Raphael *et al.* have grown both thin films and single crystals of Co_2MnSi .²¹ Although these films were found to adopt the $L1_1$ structure, there was a strong disorder between the Mn and Co sites even in the case of bulk Co_2MnSi .²² Also, Geiersbach and collaborators have grown (110) thin films of Co_2MnSi , Co_2MnGe and Co_2MnSn using a metallic seed on top of a MgO(001) substrate.²³ Finally there also exist first-principles calculations for the (001) surface of such an alloy.^{10,24}

Suits²⁵ was the first to synthesize compounds of the form Rh_2MnZ , where Z stands for Al, Ga, In, Tl, Ge, Sn, and Pb. They all crystallize in the $L2_1$ structure but compounds containing a type-II sp element show considerable disorder between the sp atom and the Mn site. They are all ferromagnets and the compounds containing Ge, Sn, and Pb have a Curie temperature above room temperature. Kanomata *et al.*²⁶ have grown crystals of the type Ru_2MnZ , where Z stands for Si, Ge, Sn, and Sb. Gotoh *et al.*²⁷ have shown that these alloys are antiferromagnets with Néel temperatures near room temperature, and Ishida *et al.*²⁸ using first-

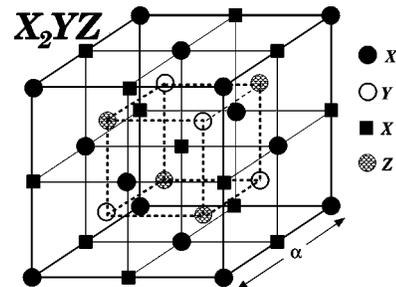


FIG. 1. Schematic representation of the $L2_1$ structure. The lattice consists of four fcc sublattices. The unit cell is that of a fcc lattice with four atoms as basis: X at (000) and $(\frac{1}{2} \frac{1}{2} \frac{1}{2})$, Y at $(\frac{1}{4} \frac{1}{4} \frac{1}{4})$, and Z at $(\frac{3}{4} \frac{3}{4} \frac{3}{4})$ in Wyckoff coordinates.

principles calculations demonstrated that the ground state is antiferromagnetic with the Mn atoms in the (111) plane being antiferromagnetically coupled to the neighboring (111) planes.

In this contribution we study the full-Heusler alloys based on Co, Fe, Ru, and Rh by extending our work on the half-Heusler alloys (see Ref. 5) and on the transition-metal monoarsenides.²⁹ To perform the calculations we have used the full-potential screened Korringa-Kohn-Rostoker (FSKKR) Green's function method³⁰ in conjunction with the local spin density approximation.³¹ The details of our calculations have been already described in Ref. 5. For all the compounds under study we have used the experimental lattice constants^{11,26} and have assumed that they are all ferromagnets. In Sec. II we present the properties of the Co_2MnZ compounds. In Sec. III we discuss the origin of the gap in these compounds and the Slater-Pauling (SP) behavior of the total moments. In Sec. IV we present our results for some other interesting systems. Finally in Sec. V we summarize our results and conclude.

II. Co_2MnZ COMPOUNDS

The first family of alloys we will be interested in are the compounds containing Co and Mn as they are the ones that have attracted most of the attention. They are all strong ferromagnets with high Curie temperatures (above 600 K) and except the Co_2MnAl they show very little disorder.¹¹ They adopt the $L2_1$ structure, which we present in Fig. 1. Each Mn or sp atom has eight Co atoms as first neighbors sitting in an octahedral symmetry position, while each Co has four Mn and four sp atoms as first neighbors and thus the symmetry of the crystal is reduced to the tetrahedral one. The Co atoms occupying the two different sublattices are chemically equivalent as the environment of the one sublattice is the same as the environment of the second one but rotated by 90° . The occupancy of two fcc sublattices by Co (or in general by X) atoms distinguish the full-Heusler alloys with $L2_1$ structure from the half-Heusler compounds with $C1_b$ structure, like, e.g., CoMnSb , where only one sublattice is occupied by Co atoms and the other one is empty. Although in the $L2_1$ structure the Co atoms are sitting on second-neighbor positions, their interaction is important to explain the magnetic properties of these compounds as we will show in the next section. In Fig. 2 we have gathered the spin-resolved total DOS for the Co_2MnAl , Co_2MnGa , Co_2MnSi , and Co_2MnGe compounds calculated using the FSKKR and in Table I the atom-projected and the total spin magnetic moment for these four compounds and for Co_2MnSn . First, as shown by photoemission experiments by Brown *et al.*³² in the case of Co_2MnSn and verified by our calculations the valence band extends 5 eV below the Fermi level and the spin-up DOS shows a large peak just below the Fermi level for these compounds. Although Ishida *et al.*¹⁵ have predicted them to be half-ferromagnets with small spin-down gaps ranging from 0.1 to 0.3 eV depending on the material, within our calculations we find that the Fermi level falls within a region of very small spin-down DOS for all these compounds. Our results agree with the calculations of Kübler

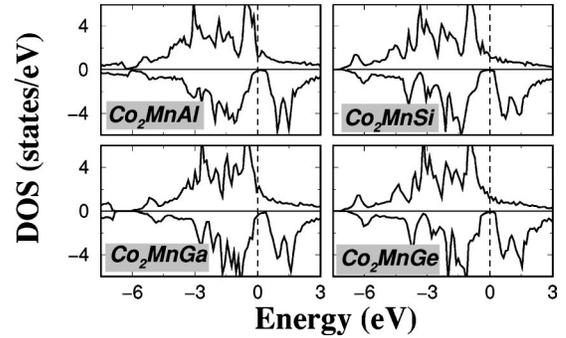


FIG. 2. Calculated spin-projected DOS for the Co_2MnZ compounds, where Z stands for Al, Ga, Si, and Ge. They all possess a finite very small spin-down DOS around the Fermi level.

*et al.*¹³ who studied the Co_2MnAl and Co_2MnSn compounds using the augmented spherical wave (ASW) method and found also a very small spin-down DOS at the Fermi level and not a real gap. The reason for this pseudogap can be found in Fig. 3 where we have drawn the band structure for the minority electrons in the case of the Co_2MnGe compound (our spin-down band structure is similar the one obtained for Co_2FeGa and Mn_2VAl in Refs. 33 and 34, respectively). We see that the Fermi level touches the highest occupied bands at the Γ point and the lowest unoccupied bands at the X point and thus the indirect gap found in the half-Heusler alloys² is practically destroyed in these materials but there is still a reasonably large direct gap at the W , K , and X points. However, we should mention that if we considerably enlarge the figure with the band structure, it can be seen that the bands do not really touch the Fermi level but there is a very small indirect gap of the order of 0.001 eV and thus the minimum of the minority unoccupied bands at X and the maximum of the occupied bands at the Γ point are not degenerated. Our calculations include relativistic effects only within the scalar-relativistic approximation; thus effects like the spin-orbit coupling can lift the band degeneracy and might even destroy the indirect gap. However, we should mention that in the case of sufficiently large band gaps like in the case of NiMnSb , the spin-orbit coupling does not destroy the half-metallicity.⁵

In the case of the half-Heusler alloys⁵ like NiMnSb the Mn spin magnetic moment is very localized due to the exclusion of the spin-down electrons at the Mn site and amounts to about $3.7\mu_B$ in the case of NiMnSb . In the case of CoMnSb the increased hybridization between the Co and

TABLE I. Calculated spin magnetic moments in μ_B using the experimental lattice constants (see Ref. 11) for the Co_2MnZ compounds, where Z stands for the sp atom.

$m^{spin}(\mu_B)$	Co	Mn	Z	Total
Co_2MnAl	0.768	2.530	-0.096	3.970
Co_2MnGa	0.688	2.775	-0.093	4.058
Co_2MnSi	1.021	2.971	-0.074	4.940
Co_2MnGe	0.981	3.040	-0.061	4.941
Co_2MnSn	0.929	3.203	-0.078	4.984

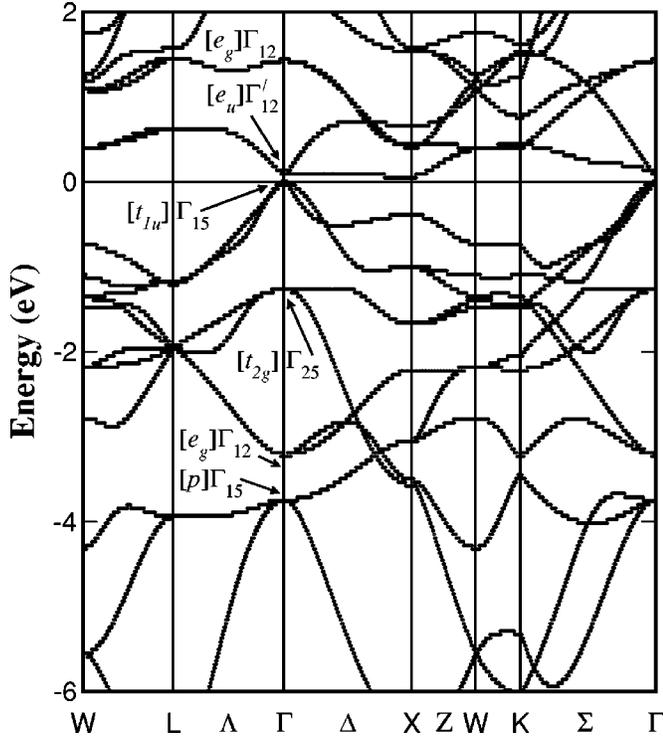


FIG. 3. Spin-down band structure of the Co_2MnGe compound. The indirect gap, present in the half-Heusler alloys, is practically destroyed. For the explanation of the different representations of the group symmetry of the Γ point look at Table II. In the brackets we present the type of orbitals transforming following each representation (see Fig. 4).

Mn spin-down electrons decreased the Mn spin moment to about $3.2\mu_B$. In the case of the full-Heusler alloys each Mn atom has eight Co atoms as first neighbors instead of four as in CoMnSb and the above hybridization is very important, decreasing even further the Mn spin moment to less than $3\mu_B$ except in the case of Co_2MnSn where it is comparable to the CoMnSb compound. The Co atoms are ferromagnetically coupled to the Mn spin moments and they possess a spin moment that varies from $\sim 0.7\mu_B$ to $1.0\mu_B$, while the sp atom has a very small negative moment which is one order of magnitude smaller than the Co moment. The negative sign of the induced sp moment characterizes most of the studied full- and half-Heusler alloys with very few exceptions.

Another important point is that in half-metallic materials like the ones studied here the total spin moment should be an integer number since both the total number of valence electrons and the number of occupied minority states are integers. However, our results in Table I do not give integer numbers for the total moments, but slight deviations of about $0.05\mu_B$. This does not arise from incorrect space integration, as it, e.g., can occur in the atomic sphere approximation. In our implementation of the full potential, the space is divided into Voronoi polyhedra,³⁰ which exactly fill up the space without any overlap, so that the space integration is performed exactly. Rather the small deviations arise from an inherent feature of the KKR Green's function method, which

due to the ℓ_{max} cutoff violates the state normalization. A proper state counting leading to integer numbers for the total charges could only be achieved if all angular momenta up to $\ell_{max} = \infty$ were included in the calculation, which is practically impossible in realistic cases. The problem can be overcome by the application of Lloyd's formula,³⁵ which contains an implicit summation over all angular momenta, thus yielding the correct total charge and moment. Since the evaluation of Lloyd's formula is a complex numerical problem, this is usually avoided arising to the above small inconsistencies.

Recently, a member of our group³⁶ has succeeded in implementing Lloyd's formula into our Green's function code, and we have tested the case of Co_2MnGe . The calculations give indeed an integer total moment of $5\mu_B$ (instead of $4.941\mu_B$ in Table I), and the (noninteger) local moments are slightly increased. Most of the charge adjustment occurs in the metallic majority band and the Fermi level is practically unchanged, situated as in Fig. 2 in the minority gap. This is also plausible from an energetic point of view; the total energy favors this position of the Fermi level. Based on this experience and on calculations as above with different ℓ_{max} cutoffs, we conclude (i) that in our calculations the correct criterium for half-metallicity is that the Fermi level is in the minority gap and (ii) that the small deviations of the total moments from integer values are insignificant.

Thus we have verified by the DOS that all compounds under study in this section are half-metals. The compounds containing Al and Ga have 28 valence electrons and the ones containing Si, Ge, and Sn 29 valence electrons. The first compounds have a total spin moment of $4\mu_B$ and the second ones of $5\mu_B$ which agree with the experimentally deduced moments of these compounds.³⁷ So it seems that the total spin moment M_t is given with respect to the total number of valence electrons, Z_t , from the simple relation $M_t = Z_t - 24$. In the following we will analyze the origin of this rule.

III. ORIGIN OF THE GAP AND SLATER-PAULING BEHAVIOR

As we mentioned above, the total spin magnetic moments of the Co_2MnZ compounds follow the $M_t = Z_t - 24$ rule. A similar relation, i.e., $M_t = Z_t - 18$, is also found for the half-Heusler compounds.^{5,38} Both state nothing more than the well-known Slater-Pauling behavior.³⁹ In such a picture the occupancy of the spin-down bands does not change and the extra or missing electrons are taken care of by the spin-up states only. The 24 means that there are 12 occupied spin-down states, as the total moment, which is the number of uncompensated spins, is given by the total number of valence electrons Z_t minus 2 times the number of minority electrons.

In Fig. 3 we present the representations of each one of the bands at the Γ point (see Table II for the different representations). First, the sp atom creates one s band and three p bands which are fully occupied. The s electrons transform following the Γ_1 representation; we do not show this band in Fig. 3 as it very low in energy and it is well separated by the other bands. The p electrons of the sp atom transform fol-

TABLE II. Representations of the real-space octahedral (O_h) symmetry group (first column). In the second column the corresponding representations of the symmetry group of the Γ point following the nomenclature introduced in Ref. 40. In the third and fourth columns we present the orbitals which transform following each one of the different representations. Notice that the whole crystal has tetrahedral T_d symmetry but the lattice consisting only of Co atoms has the O_h symmetry; T_d is a subgroup of O_h . Thus it is possible to have states located only at the Co sites, e.g., the d orbitals transforming according to the E_u representation. Also the d hybrids transforming according to the T_{1u} representation are localized at the Co atoms as there are no d states at the Mn site with the same representation. The subscripts “a” and “b” refer to orbitals at the two different Co sites in the unit cell (look Fig. 1); the 1, 2, 3, 4, and 5 refer to d orbitals of the xy , yz , zx , $3z^2-r^2$, and x^2-y^2 symmetries, respectively; the 1, 2, and 3 refer to p orbitals of the x , y , and z symmetries, respectively.

O_h	Ref. 40	Co-Co	Mn or Ge
A_{1g}	Γ_1	$s_a + s_b$	s
A_{1u}	Γ'_1	$s_a - s_b$	
E_g	Γ_{12}	$d_{ia} + d_{ib}$ [$i=4,5$]	$d_4 d_5$
E_u	Γ'_{12}	$d_{ia} - d_{ib}$ [$i=4,5$]	
T_{2g}	Γ_{25}	$p_{ia} - p_{ib}$ & $d_{ia} + d_{ib}$ [$i=1,2,3$]	$d_1 d_2 d_3$
T_{1u}	Γ_{15}	$p_{ia} + p_{ib}$ & $d_{ia} - d_{ib}$ [$i=1,2,3$]	$p_1 p_2 p_3$

lowing the Γ_{15} representation and they hybridize with p electrons of the Mn and Co atoms which transform with the same representation. As can be seen in the band structure, these bands are lower than the bands that have mainly d character but they are not well separated by them (there is a band crossing along the ΓK direction). As in the half-Heusler alloys,⁵ the four sp bands can be only partially filled by the n valence electrons of the sp atom ($n=3$ for Al and Ga or 4 for Si, Ge, and Sn), so that an additional $8-n$ d electrons are accommodated in these bands (4 d electrons in the case of Co_2MnGe or 5 d electrons for Co_2MnAl). Therefore in the Heusler alloys the effective number of d electrons (in the higher-lying d bands) can be controlled by the valence of the sp atom. This is a very unusual behavior for metallic systems, which can be used to engineer Heusler alloys with very different magnetic properties (see Sec. IV).

In the case of the half-Heusler alloys, like CoMnSb , there is only one Co atom per unit cell and its d valence electrons are hybridizing with the Mn ones creating five bonding states below the Fermi level and five antibonding ones above the Fermi level. In the full-Heusler alloys the existence of the second Co atom makes the physics of these systems more complex. As we mentioned above the whole crystal has tetrahedral symmetry (T_d). But if we neglect the Mn and the sp sites, then the Co atoms themselves sit on a cubic lattice respecting the octahedral symmetry (O_h). So there could be states obeying the O_h being localized exclusively at the Co sites; note here that the T_d is a subgroup of O_h . Thus we will take into account first the interactions between the two inequivalent Co sites and then their interaction with the Mn or the sp atom, as was also the case for the Fe_2MnZ compounds.¹⁷

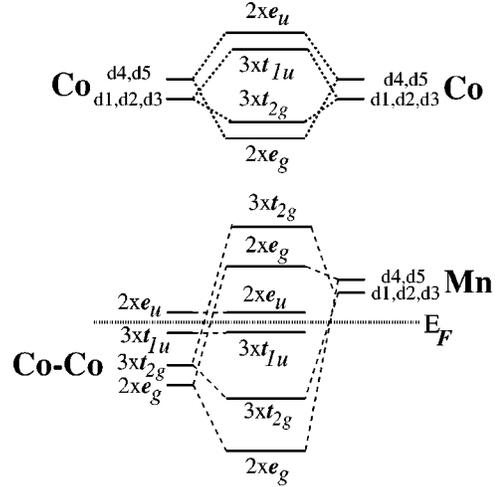


FIG. 4. Possible hybridizations between spin-down orbitals sitting at different sites in the case of the Co_2MnGe compound. To explain the properties of the full-Heusler alloys, first we consider the hybridization between the two different Co atoms and afterwards the hybridization with the Mn atom. The names of the orbitals follow the nomenclature introduced in Table II. The coefficient represents the degeneracy of each orbital.

In order to discuss the behavior of the d electrons in the full Heusler alloys we have drawn schematically in Fig. 4 the possible hybridizations between the different atoms. The $d_{1,\dots,5}$ orbitals correspond to the d_{xy} , d_{yz} , d_{zx} , $d_{3z^2-r^2}$, and $d_{x^2-y^2}$ orbitals, respectively. The symbol e_g means that the orbital transforms following the E_g representation. Note that due to symmetry, the e_g orbitals at the Co site can only couple with e_g orbitals at the other Co site or at the Mn site. The same applies for the t_{2g} orbitals. Looking at Fig. 4 we see first that when two neighboring Co atoms interact, their d_4 and d_5 orbitals form bonding e_g and antibonding e_u states; the coefficient in front of each orbital is the degeneracy of this orbital. The d_1 , d_2 , and d_3 orbitals of each Co also hybridize, creating a triple-degenerated bonding t_{2g} orbital and a triple-degenerated antibonding t_{1u} orbital.

As we show in the second part of Fig. 4, the double-degenerated e_g orbitals hybridize with the d_4 and d_5 of the Mn that transform also with the same representation. They create a double-degenerated bonding e_g state that is very low in energy and an antibonding one that is unoccupied and above the Fermi level. The $3 \times t_{2g}$ Co orbitals couple to the $d_{1,2,3}$ of the Mn and create six new orbitals, three of which are bonding and are occupied and the other three are antibonding and high in energy. Finally the $2 \times e_u$ and $3 \times t_{1u}$ Co orbitals cannot couple with any of the Mn d orbitals as there are none transforming with the u representations. The t_{1u} states are below the Fermi level and they are occupied while the e_u are just above the Fermi level. Thus in total eight minority d bands are filled and seven are empty. Our description is somewhat different from the one in Ref. 17 where it has been assumed that the orbitals just below the Fermi level are also t_{2g} and not t_{1u} as in our case. To elucidate this difference we have drawn in Fig. 5 the atomic-resolved d DOS projected on the double-degenerated and triple-

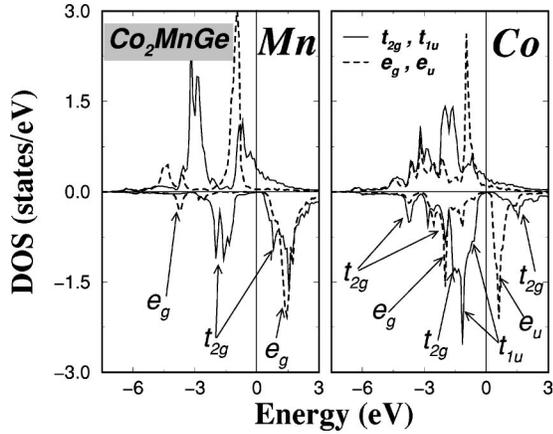


FIG. 5. Projected d DOS on the double- and triple-degenerated representations for each atom in the Co_2MnGe compound. We also give the character of each peak for the spin-down states. Notice that in the minority bands around the Fermi level there are only Co states.

degenerated representations. Although we cannot distinguish in our projection the t_{2g} from the t_{1u} and the e_g from the e_u , around the Fermi level the Mn atom presents a broad spin-down gap which is not present at the Co sites. So minority states around the gap are localized at the Co and do not couple to Mn, and the only states that have this property are the t_{1u} and e_u . Thus the peak below the Fermi level is the $3 \times t_{1u}$ state and the peak just above the Fermi level is the $2 \times e_u$ state. This also explains why the gap is small. The two cobalt atoms are second neighbors and their hybridization is not so strong and the splitting of the states is small and thus the energy distance between the t_{1u} levels and the e_u ones is small. As these states do not hybridize with the Mn states, their splitting does not change and the gap is considerably smaller than the one in the half-Heusler alloys. In the latter compounds we have only one Co atom per unit cell coupling to the Mn atom and so the t_{1u} and e_u states are absent and only the e_g and t_{2g} survive. Therefore a real gap exists in the half-Heusler alloys and the minority valence and the minority valence bands contain nine electrons: $1 \times s$, $3 \times p$, and $5 \times d$.

To summarize, in the case of the full-Heusler alloys we have eight occupied minority d states per unit cell: the double-degenerated e_g very low in energy, the triple-degenerated t_{2g} orbital, and finally the triple-degenerated t_{1u} just below the Fermi level. Thus in total we have 12 minority occupied states per unit cell, one with s character, three with p character, and eight with d character. Therefore the total moment obeys the simple rule $M_t = Z_t - 24$ as compared to $M_t = Z_t - 18$ for the half-Heusler alloys. Note here that as shown in Fig. 4 we have in total 15 spin-down d states, meaning 30 in total if we take into account both spin directions, so the states count is correct as each of the two Co atoms and the Mn one contributes totally ten d states. We can trace these states also in the spin-down band structure analyzing the character of each band at the Γ point. In Table II we have included the representations of the symmetry group of the Γ point in the reciprocal lattice using the nomenclature

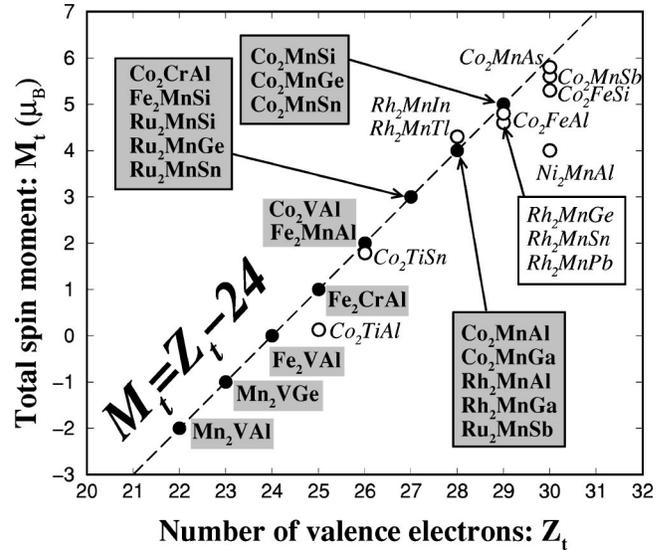


FIG. 6. Calculated total spin moments for all the studied Heusler alloys. The dashed line represents the Slater-Pauling behavior. With open circles we present the compounds deviating from the SP curve. To decide whether one alloy is half-ferromagnet or not, we have used the DOS and not the total spin-moments (see Section II).

introduced in Ref. 40. The symmetry point group of the Γ has the same symmetry operations with the O_h . First, as said above we have an s -like band not shown in the figure with a Γ_1 state at the Γ point and then we find at Γ a triple-degenerated point that has the Γ_{15} representation corresponding to the p -like orbitals. Above this point there is a double-degenerated Γ_{12} point which corresponds to the e_g orbitals while the other e_g orbitals for Co_2MnGe are found above the Fermi level and also above the unoccupied e_u orbitals that correspond to the double-degenerated point with Γ'_{12} symmetry. Finally, there are two triple-degenerated points Γ_{25} and Γ_{15} which correspond to the occupied t_{2g} and t_{1u} orbitals, respectively, while the other unoccupied t_{2g} orbitals (Γ_{25}) are high in energy and are not shown in the figure.

From the above discussion we find that in the minority band seven d states above E_F are unoccupied. Thus the largest possible moment, which a full-Heusler alloys can have, is $7 \mu_B$, since in this case all majority d states are filled. This is different from the half-Heusler compounds which have five empty d states in the minority band and therefore a maximum moment of $5 \mu_B$.

IV. OTHER FULL-HEUSLER COMPOUNDS FOLLOWING THE SP CURVE

Following the discussion of the previous section we will go on investigating other full-Heusler alloys that can follow the Slater-Pauling curve and in Fig. 6 we have plotted the total spin magnetic moments for all compounds under study as a function of the total number of valence electrons. The dashed line represents the rule $M_t = Z_t - 24$. In the following we will analyze all these results. Overall we see that many of our results coincide with the Slater-Pauling curve. Some of the Rh compounds show small deviations which are more

TABLE III. Calculated spin magnetic moments in μ_B using the experimental lattice constants (see Ref. 11) for the full-Heusler alloys containing Co, Fe, and Mn.

$m^{spin}(\mu_B)$	Co,Fe,Mn	Y	Al,Si,Ge	Total
Co ₂ TiAl	0.072	-0.013	-0.002	0.130
Co ₂ TiSn	0.911	-0.039	0.001	1.784
Co ₂ VAl	0.863	0.232	-0.033	1.926
Co ₂ CrAl	0.755	1.536	-0.091	2.955
Co ₂ MnAl	0.768	2.530	-0.096	3.970
Co ₂ FeAl	1.129	2.730	-0.099	4.890
Fe ₂ VAl	paramagnet			
Fe ₂ CrAl	-0.093	1.108	-0.011	0.910
Fe ₂ MnAl	-0.275	2.548	-0.019	1.979
Fe ₂ MnSi	0.191	2.589	-0.029	2.943
Mn ₂ VAl	-1.413	0.786	0.018	-2.021
Mn ₂ VGe	-0.750	0.476	0.021	-1.003

serious for the Co₂TiAl compound. We see that there is no compound with a total spin moment of $7\mu_B$ or even $6\mu_B$. Moreover, we found also examples of half-metallic materials with less than 24 electrons: Mn₂VGe with 23 valence electrons and Mn₂VAl with 22 valence electrons.

A. Co₂YAl and Fe₂YAl compounds

We have calculated the spin moments of the compounds Co₂YAl where Y=Ti, V, Cr, Mn, and Fe and in Table III we have gathered the atomic and total spin magnetic moments. There are experimental results only for the moment at the Co site for the Ti, V, and Cr compounds using hyperfine field measurements by Pendl *et al.*⁴¹ and by Carbonari *et al.*,⁴² which agree very well with our *ab initio* results. The compounds containing V, Cr, and Mn show a similar behavior. As we substitute Cr for Mn, which has one valence electron fewer than Mn, we depopulate one Mn spin-up state and thus the spin moment of Cr is around $1\mu_B$ smaller than the Mn one while the Co moments are practically the same for both compounds. This behavior is clearly seen in Fig. 7 where we

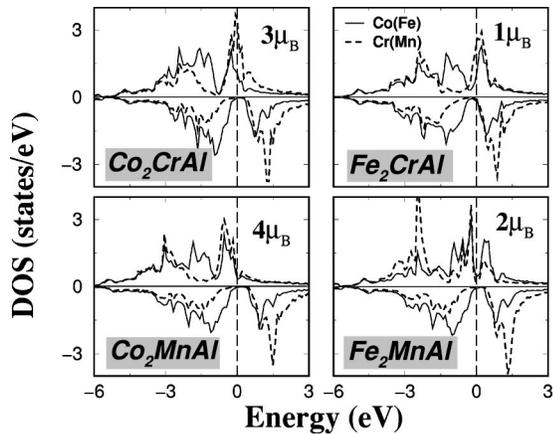


FIG. 7. Calculated atom- and spin-projected DOS for the Co(Fe)₂Mn(Cr)Al compounds. They all present a spin-down pseudogap. The numbers give the total moments.

present the atom- and spin-resolved DOS for the two compounds. The minority DOS is the same for both compounds as they follow the SP curve and this is also the case for the Co spin-up DOS. In the case of the Cr compound the Fermi level falls within a broad and large Cr spin-up peak. When we substitute Mn for Cr this peak moves lower in energy to account for the extra electron and the Fermi level is now at the right edge of the peak, but nothing else changes in the calculated DOS. Substituting V for Cr has a larger effect since also the Co spin-up DOS changes slightly and the Co magnetic moment is increased by about $0.1\mu_B$ compared to the other two compounds and V possesses a small moment of $0.2\mu_B$. This change in the behavior is due to the smaller hybridization between the Co atoms and the V compared to the Cr and Mn atoms. Although all three Co₂VAl, Co₂CrAl, and Co₂MnAl compounds are on the SP curve as can be seen in Fig. 6, this is not the case for the compounds containing Fe and Ti. If the substitution of Fe for Mn followed the same logic as the one of Cr for Mn, then the Fe moment should be around $3.5\mu_B$ which is a very large moment for the Fe site. Therefore it is energetically more favorable for the system that also the Co moment is increased, as was also the case for the other systems with 29 electrons like Co₂MnSi, but while the latter one makes it to $5\mu_B$, Co₂FeAl reaches a value of $4.9\mu_B$. A similar behavior was seen also in the case of the isoelectronic Co₂FeGa compound, but the total spin moment was slightly larger than $5\mu_B$.³³ In the case of Co₂TiAl, it is energetically more favorable to have a weak ferromagnet than an integer moment of $1\mu_B$ as it is very difficult to magnetize the Ti atom. Even in the case of the Co₂TiSn the calculated total spin magnetic moment of $1.78\mu_B$ (compared to the experimental value of $1.96\mu_B$) (Ref. 43) arises only from the Co atoms as was also shown experimentally by Pendl *et al.*,⁴¹ and the Ti atom is practically paramagnetic and the latter compound fails to follow the SP curve.

As a second family of materials we have calculated also the compounds containing Fe and we present their total spin moments also in Table III. Fe₂VAl has in total 24 valence electrons and is a semimetal, i.e., paramagnetic with a very small DOS at the Fermi level, as is already known experimentally.⁴⁴ All studied Fe compounds follow the SP behavior as can be seen in Fig. 6. In the case of the Fe₂CrAl and Fe₂MnAl compounds the Cr and Mn atoms have spin moments comparable to the Co compounds and similar DOS as can be seen in Fig. 7. In order to follow the SP curve the Fe in Fe₂CrAl is practically paramagnetic while in Fe₂MnAl it has a small negative moment. When we substitute Si for Al in Fe₂MnAl, the extra electron exclusively populates Fe spin-up states and the spin moment of each Fe atom is increased by $0.5\mu_B$ contrary to the corresponding Co compounds where also the Mn spin moment was considerably increased.

Finally we calculated as a test Mn₂VAl and Mn₂VGe which have 22 and 23 valence electrons, respectively, to see if we can reproduce the SP behavior not only for compounds with more than 24, but also for compounds with fewer than 24 electrons. As we have already shown Fe₂VAl is paramagnetic and Co₂VAl, which has two electrons more, has a spin moment of $2\mu_B$. Here Mn₂VAl has two valence electrons

TABLE IV. Calculated atom-resolved and total spin magnetic moments in μ_B using the experimental lattice constants for the full-Heusler alloys containing Rh and Ru (see Ref. 11 for the lattice constants of the Rh compounds and Ref. 26 for the Ru compounds).

$m^{spin}(\mu_B)$	Ru, Rh	Mn	Z	Total
Ru ₂ MnSi	0.028	2.868	0.025	2.948
Ru ₂ MnGe	0.002	2.952	0.021	2.977
Ru ₂ MnSn	-0.051	3.137	-0.001	3.034
Ru ₂ MnSb	0.222	3.495	0.018	3.957
Rh ₂ MnAl	0.328	3.388	-0.041	4.004
Rh ₂ MnGa	0.312	3.461	-0.033	4.052
Rh ₂ MnIn	0.269	3.720	-0.034	4.223
Rh ₂ MnTl	0.266	3.765	-0.027	4.270
Rh ₂ MnGe	0.421	3.672	0.011	4.525
Rh ₂ MnSn	0.393	3.831	-0.010	4.607
Rh ₂ MnPb	0.383	3.888	-0.009	4.644

fewer than Fe₂VAl and as we show in Table III its total spin moment is $-2\mu_B$, in agreement with previous *ab initio* results,³⁴ and thus it follows the SP behavior. To our knowledge there is no compound with 23 valence electrons, which has been studied experimentally, so we decided to calculate Mn₂VGe using the lattice constant of Mn₂VAl. We have chosen this compound, because as can be seen in Ref. 11 the compounds containing Al and Ge have practically the same lattice constants. We found that adding one electron to Mn₂VAl results in a decrease of the absolute value of both the Mn and V spin moments (note that V and Mn are antiferromagnetically coupled) so that the resulting Mn₂VGe total spin magnetic moment is $-1\mu_B$ following the SP curve as can be also seen in Fig. 6.

B. Ru and Rh compounds

To investigate further the Slater-Pauling behavior of the full-Heusler alloys we studied the ones containing a *4d* transition-metal atom. As we have already mentioned in Sec. I the Ru compounds are antiferromagnets with Néel temperatures that reach room temperature. We have calculated their properties assuming that they are ferromagnets and present the calculated spin-magnetic moments in Table IV. The Ru₂MnSi, Ru₂MnGe, and Ru₂MnSn, have a total spin magnetic moment of $3\mu_B$ and Ru₂MnSb a moment of $4\mu_B$ following the rule for the magnetic moments that we have already shown for the Co and Fe compounds and thus the Fermi level falls within the pseudogap contrary to the calculations in Ref. 45 where the Fermi level was above the gap. In the case of the alloys with Si, Ge, and Sn the Ru atom has a practically zero spin moment and the total moment is carried by the Mn atoms. In Fig. 8 we have drawn the atomic and spin DOS for the Ru₂MnSi compound compared to the isoelectronic Fe₂MnSi compound. We see clearly from the DOS that the hybridization between the Mn and Ru spin-down states is smaller than in the case of the Fe compound, resulting in a larger Mn spin moment. Although Ru has a practically zero spin moment, we see that the Fermi level

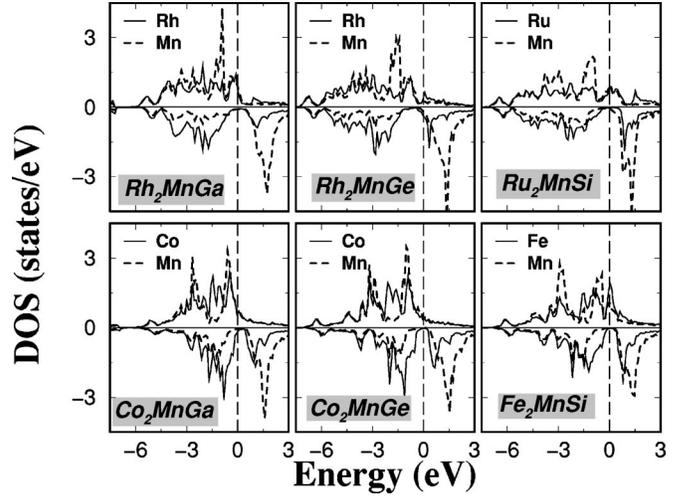


FIG. 8. Calculated atom- and spin-projected DOS for some of the Heusler alloys containing Ru and Rh compared to the alloys containing Fe or Co that are isoelectronic to Ru and Rh, respectively. In the case of the Rh or Ru compounds, the hybridization with the spin-down Mn states is smaller resulting in larger Mn spin moments (see Table IV).

falls within a broad peak of spin-up DOS. In the case of Ru₂MnSb this peak is completely occupied, resulting in an important induced spin moment at the Ru site that couples ferromagnetically to the Mn one.

The next family of compounds that we have studied are the ones containing Rh and Mn.^{25,46} In Table IV we present the calculated spin magnetic moments. As can be seen in Fig. 8 the hybridization between the Rh and the Mn spin-down states is smaller than in the case of the isoelectronic Co compounds; i.e., there are Mn states in the Co compound that become Rh states in the Rh compound, thus leading to an increase of the Mn moment and a decrease of the Rh moment compared to the Co spin moment. This phenomenon is quite intense as the Mn moment increases in all cases more than $0.6\mu_B$. From the studied compounds only Rh₂MnAl and Rh₂MnGa are exactly on the SP curve presented in Fig. 6. The Rh₂MnIn and Rh₂MnTl that are isoelectronic to the two previous compounds have a total spin moment of around $4.2\mu_B$ – $4.3\mu_B$, thus the Fermi level is slightly below the pseudogap in these compounds. In the case of Rh₂MnGe, Rh₂MnSn, and Rh₂MnPb, which possess 29 valence electrons, the total spin moment is around $4.6\mu_B$ slightly smaller than the ideal $5\mu_B$ and the Fermi level is slightly above the pseudogap. This is probably due to the considerably larger lattice constant of the Rh compounds with respect to the isoelectronic Co ones. But in general, as can be seen also in Fig. 6, where we summarize all our results, all the compounds are not very far from the SP curve and the deviations are small.

C. Compounds with 30 valence electrons

As stated in Sec. III the maximal moment of a full-Heusler alloy is $7\mu_B$ and should occur when all 15 majority *d* states are occupied. Analogously for a half-Heusler alloy

TABLE V. Calculated spin moments for full-Heusler alloys containing 30 valence electrons per unit cell. The experimental lattice parameters were taken from Ref. 11.

$m^{spin}(\mu_B)$	X	Y	Z	Total
Ni ₂ MnAl	0.364	3.359	-0.062	3.973
Co ₂ FeSi	1.271	2.756	-0.031	5.268
Co ₂ MnSb	1.113	3.401	-0.007	5.620
Co ₂ MnAs	1.219	3.309	0.035	5.782

the maximal moment is $5\mu_B$. However, this limit is difficult to achieve, since due to the hybridization of the d states with empty sp states of the transition-metal atoms (sites X and Y in Fig. 1), the d intensity is transferred into states high above E_F , which are very difficult to occupy. While we could identify in a recent paper on half-Heusler alloys (Ref. 5) systems with a moment of nearly $5\mu_B$, the hybridization is much stronger in the full-Heusler alloys so that a total moment of $7\mu_B$ seems to be impossible. Therefore we restrict our search to possible systems with $6\mu_B$, i.e., systems with 30 valence electrons. We have studied some of the possible candidates and we present our results in Table V. One obvious way to reach the 30 electrons is to substitute, e.g., in Co₂MnAl, Co by Ni, but Ni is practically paramagnetic and cannot carry a large spin moment and thus the total spin magnetic moment of Ni₂MnAl is only $4\mu_B$ far away from the ideal $6\mu_B$. The second way to achieve 30 electrons is to use Fe at the Y site as is the case for the Co₂FeSi compound. Already Co₂FeAl was not reaching the $5\mu_B$ and adding one more electron cannot increase the total spin moment by more than $1\mu_B$. Although the Co moment reaches $1.3\mu_B$, the Fe moment stays unchanged and the total spin moment is increased only by $\sim 0.4\mu_B$, reaching the $5.3\mu_B$ instead of the ideal $6\mu_B$.

Our last test cases are the Co₂MnSb and Co₂MnAs compounds. We have calculated Co₂MnSb using the lattice constant of Co_{1.5}MnSb as Co₂MnSb does not really exist. Adding Co to Co_{1.5}MnSb results in the creation of a Co-rich phase. Co₂MnSn has a total spin moment of $5\mu_B$. The additional electron increases both the Co and Mn spin moments and the total moment is now $5.6\mu_B$. To our knowledge there is no experimental work on Co₂MnAs and we have calculated it using the lattice constant of Co₂MnGe. This lattice constant should be very close to the real one as also substituting Ga for Ge only marginally changes it. As shown in Table V the calculated total spin moment is $5.8\mu_B$. But for both compounds, if we increase their lattice constant by 4%,

the Fermi level moves deeper in energy, as was the case also for the half-Heusler alloys,⁵ and now it falls within the pseudogap and the total spin moment for both of them reaches the ideal value of $6\mu_B$. So if both Co₂MnSb and Co₂MnAs can be grown on top of a substrate with the appropriate lattice constant using a technique like molecular beam epitaxy, it is possible to get a material with a total spin moment of $6\mu_B$ where the Fermi level will be within the pseudogap. In such a case, of course, there is the possibility that the lattice parameter along the growth axis is contracted to account for the large in-plane lattice parameter, which can lead to a change of the total spin moment.

V. CONCLUSIONS

Using the full-potential screened Korringa-Kohn-Rostoker method we studied the full-Heusler alloys containing Co, Fe, Rh, and Ru. We have shown using the scalar-relativistic approximation that for all these compounds the top edge of the highest occupied spin-down band and the bottom edge of the lowest unoccupied spin-down band touch the Fermi level, practically destroying the indirect gap. These compounds show a Slater-Pauling behavior and the total spin-magnetic moment per unit cell (M_t) scales with the total number of valence electrons (Z_t) following the rule $M_t = Z_t - 24$. The Co-Co hybridization is primordial to explain why the spin-down band contains exactly 12 electrons and why only a tiny gap exists in these compounds. Finally we have shown that it is possible to find the Slater-Pauling behavior even for materials with fewer than 24 valence electrons like Mn₂VAl and Mn₂VGe, and that the compounds with 30 valence electrons are unlikely to achieve a total spin moment of $6\mu_B$.

Note added in proof. After acceptance we became aware of a paper⁴⁷ where the pressure dependent properties of the Co₂MnX ($X=Si, Ge, Sn$) compounds have been studied using the full-potential linearized augmented plane wave (FLAPW) method.

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